Pentafluorophenoxy Boron Subphthalocyanine (F₅BsubPc) as a Multifunctional Material for Organic Photovoltaics

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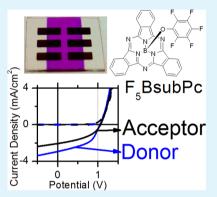
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Supporting Information

ABSTRACT: We have demonstrated that pentafluoro phenoxy boron subphthalocyanine ($F_5BsubPc$) can function as either an electron donor or an electron acceptor layer in a planar heterojunction organic photovolatic (PHJ OPV) cell. $F_5BsubPc$ was incorporated into devices with the configurations ITO/MoO₃/ $F_5BsubPc/C_{60}/BCP/Al$ ($F_5BsubPc$ used as an electron-donor/hole-transport layer) and ITO/MoO₃/Cl-BsubPc/ $F_5BsubPc/BCP/Al$ ($F_5BsubPc/BCP/Al$ ($F_5BsubPc/BCP/Al$ ($F_5BsubPc/BCP/Al$ ($F_5BsubPc$) was an electron-acceptor/electron-transport layer). Each unoptimized device displayed open-circuit photopotentials (V_{oc}) close to or in excess of 1 V and respectrable power conversion efficiencies. Ultraviolet photoelectron spectroscopy (UPS) was used to characterize the band-edge offset energies at the donor/acceptor junctions. HOMO and LUMO energy level offsets for the $F_5BsubPc/C_{60}$ heterojunction were determined to be ca. 0.6 eV and ca. 0.7 eV, respectively. Such offsets are clearly large enough to produce rectifying J/V responses, efficient exciton dissociation, and photocurrent production at the interface. For the Cl-BsubPc/ $F_5BsubPc$ heterojunction, the estimated offset energies were found to be ca.



0.1 eV. However, reasonable photovoltaic activity was observed, with photocurrent production coming from both BsubPc species layers. Incident and absorbed photon power conversion efficiencies (IPCE and APCE) showed that photocurrent production qualitatively tracked the absorbance spectra of the donor/acceptor heterojunctions, with some additional photocurrent activity on the low energy side of the absorbance band. We suggest that photocurrent production at higher wavelengths may be a result of charge-transfer species at the donor/acceptor interface. Cascade photovoltaics were also fabricated to expand on the understanding of the role of $F_5BsubPc$ in such device architectures.

KEYWORDS: photovoltaic, heterojunctions, solar, cells, boron, subphthalocyanines

INTRODUCTION

The exploration of new classes of organic semiconductor materials continues to be of interest especially the exploration of new materials that can function as active layers in multiple device configurations such as organic light emitting diodes (OLEDs), organic photovoltaics (OPVs), and organic field effect transistors (OFETs), especially if there is a possibility that they can transport both holes and electrons with reasonable charge carrier mobilities.^{1–3}

Boron subphthalocyanines (BsubPcs) are a class of strong magenta colored $14-\pi$ electron aromatic materials that have recently shown to be organic semiconductors in OPVs,^{4–6} OLEDs,^{7–9} and OFETs.¹⁰ In these applications, BsubPcs were used as either an electron donor (hole transport material)^{4,10} or an electron acceptor (electron transport material).⁶ Included in this list is our demonstration that a series of fluorinated phenoxy boron subphthalocyanine, exemplified by pentafluoro phenoxy boron subphthalocyanine (F₅BsubPc), can function as both electron-transporting and electroluminescent materials in

OLEDs.^{7,8} More recently we have also measured the field dependent electron mobilities of the same series of fluorinated phenoxy boron subphthalocyanines and have shown their electron mobilities to be comparable to those seen for several other known electron-transporting materials (mobilities up to ca. 10^{-4} cm²/V·s were observed).¹¹ We have shown that a novel series of phthalimido-BsubPcs can be used as electron-transporting emitters in OLEDs. The stable one-electron reduction and one-electron oxidation processes observed in solution for the phthalimido-BsubPcs suggest that they may be applicable as both electron- and hole-transporting materials.¹² Jones and coworkers have also recently shown that hole and electron transport is possible through a layer of chloro boron subphthalocyanine (Cl-BsubPc) in an OPVs.¹³

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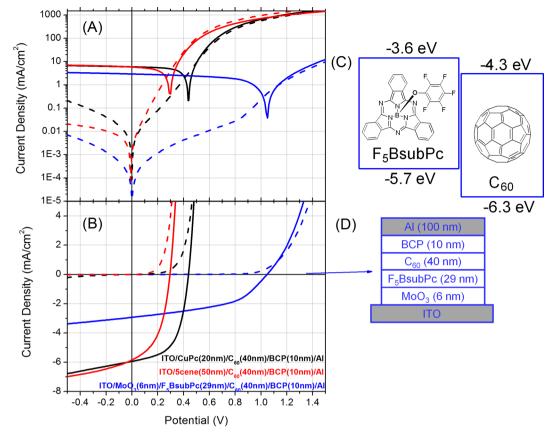


Figure 1. Log (A) and linear (B) dark and illuminated J/V plots for OPV devices fabricated with $F_5BsubPc$ (blue line), CuPc (black line), or PEN (red line) as donor materials and C_{60} as an acceptor. (C) Schematic of the molecular structure of $F_5BsubPc$ and its frontier molecular energy level alignment with C_{60} . (D) Device configuration of the $F_5BsubPc/C_{60}$ OPV.

Recently, in this journal, our group has extensively reviewed the application of BsubPcs in organic electronic devices.¹⁴ In that review we point out that the majority of the research concerning BsubPcs in organic electronic devices has been centered on the prototypical BsubPc: Cl-BsubPc. At the end of the review, we briefly mention preliminary results wherein we illustrate that $F_5BsubPc$, like Cl-BsubPc,⁶ can function as either a electron-donor/hole-transporting layer or an electron-acceptor/electron-transporting layer in an OPV. We then go on to make the case for the expanded exploration of phenoxy-BsubPcs and other axially modified BsubPcs in organic electronic devices owing to the wide structure variation possible and the potential for systematically tuning the physical and electrophysical properties through synthetic variation.

Herein we present a comprehensive study of $F_5BsubPc$ in planar heterojunction (PHJ) OPVs, $F_5BsubPc$ being a representative from the more general class of materials, phenoxy-BsubPcs. We show that $F_5BsubPc$ can function as an electron-donor material in a planar heterojunction (PHJ) OPV using C_{60} as the electron-accepting material. We also show that $F_5BsubPc$ can function as an electron-acceptor material using Cl-BsubPc as the electron-donating material in an "all-BsubPc" PHJ OPV. Ultraviolet photoelectron spectroscopy (UPS) was used to estimate the ionization potentials for each layer in the PHJ OPV configurations, allowing for the estimation of the HOMO energy levels (and as a derivative thereof the LUMO energy levels) of the active materials and an estimation of the resulting impact of heterojunction formation between $F_5BsubPc$ and either C_{60} or Cl-BsubPc on the energy levels.

RESULTS AND DISCUSSION

OPVs Based on F₅BsubPc/C₆₀ Heterojunctions. BsubPcs are most commonly used in OPVs as electron-donor materials, the most common of which is Cl-BsubPc.^{4,14} OPVs based on the pairing of Cl-BsubPc and C₆₀ have shown good performance with the following device architecture: ITO/ MoO₃ (5 nm)/Cl-BsubPc (10 nm)/C₆₀ (33 nm)/BCP (8 nm)/Al (100 nm).^{15,16} Therefore, F₅BsubPc was first evaluated as an alternative electron-donor material to Cl-BsubPc. F₅BsubPc/C₆₀ devices were constructed as follows: ITO/ MoO₃/F₅BsubPc/C₆₀/BCP/Al (Figure 1D). MoO₃ was used as an interlayer because it has been shown to enhance the efficiency of hole harvesting for a variety of electron donor materials.¹⁶⁻¹⁹ The F₅BsubPc layer was varied in thickness from 17 to 37 nm, while C_{60} thickness was held constant at 40 nm. The devices consistently produced a large $V_{\rm oc}$ of ${\sim}1~{\rm V}$ (Table 1), larger than comparative devices constructed using electron donors such as copper phthalocyanine (CuPc) and pentacene (PEN, Figure 1, Table 1).⁴ The F₅BsubPc layer thickness had a minimal impact on open circuit potential (V_{oct} varying by ± 0.05 V) and fill factors (FF, varying between 0.4 to 0.5) across the thickness range (Table 1). At 37 nm, the FF rapidly dropped to 33% and the JV curve resembled an s-kink, which we believe is caused by unbalanced charge transport at this thickness, leading to excessive recombination losses.²⁰⁻²² The power conversion efficiency was mainly dictated by the changes in current produced as a result of the change in exciton dissociation efficiency as the product of exciton diffusion length, $L_{\rm D}$, and absorptivity, α , exceeded the optimal $\alpha L_{\rm D}$

Table 1. Summary of Device Optimization by Varying the $F_5BsubPc$ Thickness in the Following Device Configuration: ITO/MoO₃ (6 nm)/F₅BsubPc/C₆₀ (40 nm)/BCP (10 nm)/Al (100 nm)^{*a*}

F ₅ BsubPc thickness (nm)	$V_{\rm oc}$ (V)	J _{sc} (mA/cm ²)	FF (%)	η (%)
17	1.07 ± 0.01	0.88 ± 0.16	48.6 ± 1.9	0.46 ± 0.10
27	1.02 ± 0.03	2.03 ± 0.39	46.2 ± 3.9	0.96 ± 0.22
29	1.05 ± 0.02	2.96 ± 0.59	48.2 ± 1.7	1.49 ± 0.27
32	1.05 ± 0.00	2.85 ± 0.63	43.9 ± 1.0	1.32 ± 0.27
37	1.03 ± 0.03	2.32 ± 0.32	33.4 ± 0.9	0.80 ± 0.16
reference CuPc (20 nm)	0.44 ± 0.02	5.94 ± 0.79	56.9 ± 2.5	1.49 ± 0.26
reference PEN (50 nm)	0.30 ± 0.01	5.84 ± 1.59	46.9 ± 1.6	0.82 ± 0.29

^aStandard deviations for an average of at least eight 0.019 cm² devices are reported for each performance parameter. Two reference devices are shown for comparisons: ITO/CuPc (20 nm)/C₆₀ (40 nm)/BCP (10 nm)/Al (100 nm) and ITO/PEN (50 nm)/C₆₀ (40 nm)/BCP (10 nm)/Al (100 nm).

product.^{23,24} It is also notable that the currents at zero bias are significantly lower than for the CuPc and PEN based OPVs, consistent with the idea that reverse saturation currents in the conventional Shockley equation describing the JV response of an OPV are dictated by the energy barrier between the transport HOMO of the donor and the transport LUMO of the

acceptor and energy that is significantly higher for the BsubPcs.²⁵ Overall, the best performing device was found to be ITO/MoO₃ (6 nm)/F₅BsubPc (29 nm)/C₆₀ (40 nm)/BCP (10 nm)/Al (100 nm) with an average power conversion efficiency of 1.5% (champion devices were observed with efficiencies as high as 2% (Figure 1A,B; Table 1)).

F_sBsubPc can clearly function as an electron donor when paired with C₆₀; however, F₅BsubPc also functions as an electron transporter.^{7,11} Organic materials have been shown to form Schottky barriers with MoO₃ interfaces, producing photocurrent (organic material acting as electron acceptor).² Therefore, we were interested in locating the rectifying interface for exciton dissociation in the $F_5BsubPc/C_{60}$ device(s) to prove the function of F5BsubPc as an electron donor material. We sequentially "deconstructed" the $F_5BsubPc/C_{60}$ device by making a series of new devices by first omitting the C_{60} layer and then the MoO₃ layer (Figure 2). We observed that the $MoO_3/F_5BsubPc$ interface(s) do not adequately rectify and produced only a minimal photocurrent. A large $V_{oc} > 1$ V was however produced due to the large 2.1 eV bandgap of $F_5BsubPc$ (Figure 2, blue line). This is similar to the large V_{oc} in MoO_3/C_{60} Schottky devices previously reported.²⁶ In contrast the $F_5BsubPc/C_{60}$ device (without MoO₃, Figure 2, red line) produced similar J_{sc} current to its reference device (Figure 2, black line), albeit with an s-kink due to charge extraction limitations presumably caused by poor alignment of the ITO 4.7 eV work function with the low lying HOMO of F₅BsubPc (IP of 5.7 eV) in the absence of MoO₃. Thus, we can conclude that the $F_5BsubPc/C_{60}$ interface is indeed the rectifying

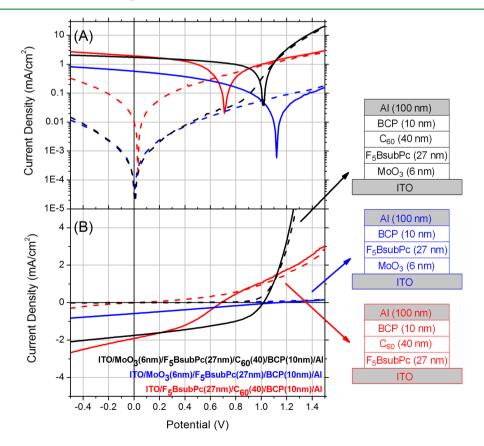


Figure 2. Log (A) and linear (B) dark and illuminated J/V plots for OPV devices: (black line) ITO/MoO₃ (6 nm)/F₅BsubPc (27 nm)/C₆₀ (40 nm)/BCP (10 nm)/Al (100 nm), (red line) ITO/F₅BsubPc (27 nm)/C₆₀ (40 nm)/BCP (10 nm)/Al (100 nm), and (blue line) ITO/MoO₃ (6 nm)/F₅BsubPc (27 nm)/Al (100 nm).

interface where photocurrent is generated in the ITO/MoO₃/ $F_5BsubPc/C_{60}/BCP/Al$ device(s), observations that are consistent with the UPS data outlined in the subsequent sections of this paper.

An IPCE plot for a representative ITO/MoO₃/ F_5 BsubPc/ C₆₀/BCP/Al large area device is shown in Figure 3. A sharp

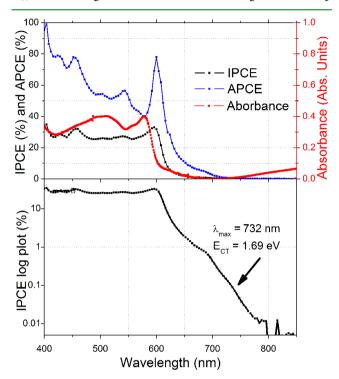


Figure 3. (top) IPCE (black line), APCE (blue line), and absorbance (red line) of an ITO/MoO₃ (6 nm)/F₅BsubPc (29 nm)/C₆₀ (40 nm)/ BCP (10 nm)/Al (100 nm) large area device. The absorbance spectrum coincides with what is expected of a combination of F₅BsubPc and C₆₀, the IPCE response peaks ca. 10–20 nm to the red of the main absorbance peak, suggesting that exciton dissociation is favored by excitation of a minority species at the F₅BsubPc/C₆₀ heterojunction. (bottom) A plot of the log of IPCE versus wavelength for the same device.

increase in the IPCE spectra at ~600 nm is seen with an IPCE response of 33% corresponding to the contribution to the photocurrent production from the absorption of light from $F_5BsubPc$. The $F_5BsubPc/C_{60}$ device also shows a significant photoresponse from the absorption of light by C_{60} between 400 and 500 nm (Figure 3). The absorption is exclusively from C_{60} , since F₅BsubPc has an absorption minimum over that range. The devices maintain ~30% IPCE throughout this absorption range. The J_{sc} estimated from the integration of the IPCE spectrum, normalized to the AM 1.5G spectral response, and the excitation lamp output,²⁷ was 1.81 mÅ/cm², which is lower than the measured value (3.1 mA/cm²). For comparison, Cl-BsubPc $(13 \text{ nm})/C_{60}$ (33 nm) devices are reported to have an IPCE of ~45% at the Cl-BsubPc absorption maximum.¹⁵ Photocurrent from a possible charge transfer (CT) complex between F₅BsubPc and C₆₀ can be observed in the IPCE spectrum, at a wavelength of 732 nm, which is consistent with absorption from a CT state at 1.69 eV (Figure 3, bottom). Similar CT states have been previously observed in Cl-BsubPc/ C₆₀ planar or graded heterojunctions with IPCE contributions starting at 600 nm and tailing to 750 nm.^{28,29} The energy of charge transfer states in OPV devices has recently been

correlated to measured $V_{\rm oc}$.³⁰ Applying this concept to the $F_5BsubPc/C_{60}$ devices predicts an $V_{\rm oc}$ of 1.09 \pm 0.07 V and in good agreement with our experimental measurements.

OPVs Based on Cl-BsubPc/F₅BsubPc Heterojunctions. In our previous work detailing the use of F₅subPc in OLEDs and in single carrier devices, we demonstrated that F₅BsubPc can act as an electron transporting and electroluminescent material.^{7,8,11} In the preceding section, we outlined the pairing of F₅BsubPc with C₆₀ wherein its role was as a hole transporter and light absorber. The frontier molecular orbitals of F₅BsubPc also suggest that it would be an effective material in OPVs as an electron acceptor if electron donor materials are selected with an appropriate energetic offset. For example, our UPS data outlined in the following sections of this paper suggest that F₅BsubPc can used as an acceptor material when paired with Cl-BsubPc (5.6 eV),⁴ albeit with a relatively small band edge offset of ~0.1 eV. There is literature precedence of "all-BsubPc" devices producing V_{oc} near 1 V with only 0.1–0.3 eV offset in their IP (0.2-0.6 eV offset in electron affinity values) and producing high power conversion efficiencies (2.7% and 4%).^{6,31} We also selected pentacene (PEN) as an alternative electron-donating material for pairing with F₅BsubPc because of its lower IP (4.8 eV) and its significant HOMO offset versus $F_5BsubPc$ (Figure 4).^{32,33}

Devices with the configurations Al (100 nm)/BCP (8 nm)/ $F_{s}BsubPc$ (x nm)/Cl-BsubPc (10 nm)/MoO₃ (8 nm)/ITO and Al (100 nm)/BCP (8 nm)/F₅BsubPc (x nm)/PEN (50 nm)/ITO (Figure 4D) were thus constructed and evaluated as PHJ OPVs. Reasonable rectification was observed in the dark for both device configurations (Figure 4). The F₅BsubPc layer thickness was varied while the thicknesses of the PEN and Cl-BsubPc layers were held constant at 50 and 10 nm, respectively (according to best practices for each outlined in the literature).^{4,34} Using $F_5BsubPc$ as the electron transport layer, $V_{\rm oc}$ values of 0.95 and 1.05 V were observed for PHJ OPVs with PEN and Cl-BsubPc electron donors, respectively (Figure 4). The best thickness for the F₅BsubPc layer in the F₅BsubPc/Cl-BsubPc devices over the range tested was 40 nm, resulting in a J_{sc} of 2.1 mA/cm². For the PEN/F₅BsubPc devices the best thickness over the range tested was found to be 51 nm producing a J_{sc} of 0.73 mA/cm². The fill factors for each device were comparable at 43.2% and 41.8%, respectively and the best power conversion efficiencies for Cl-BsubPc/F5BsubPc and PEN/F₅BsubPc devices were 0.94% and 0.28%, respectively (Table 2). The performance of larger area devices is summarized in the Supporting Information accompanying this paper.

In the case of the "all-BsubPc" Cl-BsubPc/F₅BsubPc device, IPCE measurements showed that absorption of wavelengths between 450 and 650 nm contributed to the photoresponse of the device (Figure 5). This narrow spectral response is not surprising and is a result of the near identical absorption profiles of the two active materials (Cl-BsubPc and F₅BsubPc). The highest photoresponse is observed at 596 nm with an IPCE of 47%. A second peak in the IPCE spectra is observed at 548 nm similar to the absorbance shoulder present for the BsubPcs. Charge transfer states are also observed at the onset of the IPCE plot (Figure 5, bottom). Here a peak is observed at 720 nm representing absorbance and photocurrent production from a CT state at 1.72 eV. Calculating the predicted V_{oc} using the correlation developed by Graham³⁰ suggests a value of 1.12 \pm 0.07 V, once again in good agreement with our experimentally measured V_{oc} .

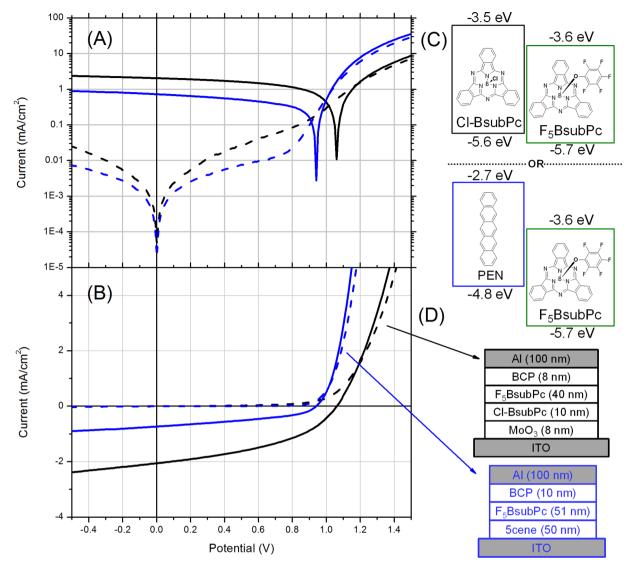


Figure 4. Log (A) and linear (B) dark and illuminated J/V plots for OPVs using $F_5BsubPc$ as an acceptor and PEN (blue line) or Cl-BsubPc (black line) as donor materials. (C) Schematics of the molecular structure of $F_5BsubPc$ and its frontier molecular energy level alignment with PEN and Cl-BsubPc. (D) The structure of the Cl-BsubPc/ $F_5BsubPc$ and the PEN/ $F_5BsubPc$ devices.

Table 2. Summary of Device Optimization by Varying the $F_5BsubPc$ Thickness in the Following Device Configurations: ITO/PEN (50 nm)/ $F_5BsubPc/BCP$ (10 nm)/Al (100 nm) and ITO/MoO₃ (8 nm)/Cl-BsubPc (10 nm)/ $F_5BsubPc/BCP$ (10 nm)/Al (100 nm)^{*a*}

F ₅ BsubPc thickness (nm)	$V_{\rm oc}$ (V) $J_{\rm sc}$ (mA/cm ²)		FF (%)	η (%)					
PEN(50 nm)/F ₅ BsubPc									
29	0.82 ± 0.04	0.46 ± 0.09	40.8 ± 6.8	0.15 ± 0.01					
38	0.89 ± 0.01	0.57 ± 0.09	42.7 ± 0.7	0.22 ± 0.03					
51	0.94 ± 0.02	0.73 ± 0.13	41.4 ± 3.5	0.28 ± 0.05					
Cl-BsubPc(10 nm)/F ₅ BsubPc									
20	0.98 ± 0.02	1.72 ± 0.52	45.1 ± 1.5	0.76 ± 0.24					
30	0.98 ± 0.02	2.01 ± 0.66	43.7 ± 2.1	0.86 ± 0.28					
40	1.06 ± 0.03	2.06 ± 0.71	43.2 ± 0.9	0.94 ± 0.33					
$^a{\rm Standard}$ deviations for an average of at least eight 0.019 ${\rm cm}^2$ devices are reported for each performance parameter.									

By integrating the IPCE spectra and normalizing it to the spectral output of our AM1.5G illumination source,²⁷ we

estimate a device $J_{\rm sc}$ of 2.38 mA/cm², which is in good agreement with experimental observations (2.19 mA/cm²). By comparison to the F₅BsubPc/C₆₀ device discussed above, an improvement from 33% to 47% in IPCE at 600 nm indicates an additional contribution from the Cl-BsubPc in the photoresponse. The IPCE performance of Cl-BsubPc/C₆₀ devices in literature^{15,29} and Cl-BsubPc/F₅BsubPc devices outlined herein are similar to a photoresponse of ~45% at 600 nm.

Finally, to probe the Cl-BsubPc/F₅BsubPc interface as the rectifying interface, we also deconstructed the ITO/MoO₃/Cl-BsubPc/F₅BsubPc/BCP/Al device. Schottky junction devices that lacked the F₅BsubPc layer were constructed (ITO/MoO₃/Cl-BsubPc/F₅BsubPc/BCP/Al). These Schottky junction devices lacked rectification (Figure S4, Supporting Information) similar to previous reports,³⁵ and therefore we can conclude that the Cl-BsubPc/F₅BsubPc must be acting as an electron acceptor, an observation that is consistent with the UPS data outlined below.

Band-Edge Energies for $F_5BsubPc/C_{60}$ and Cl-BsubPc/ $F_5BsubPc$ Heterojunctions. We then moved to establish the

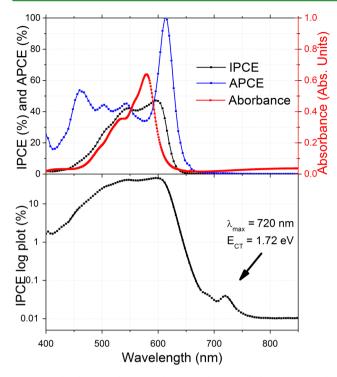


Figure 5. (top) IPCE (black line), APCE (blue line) (arbitrary units), and absorbance (red line) of the ITO/MoO₃ (8 nm)/Cl-BsubPc (10 nm)/F₃BsubPc (40 nm)/BCP (8 nm)/Al (100 nm) large area device. As for the F₃BsubPc/C₆₀ OPVs, the absorbance spectrum is a composite of absorbance from both BsubPcs; the IPCE spectra suggest photocurrent production from minority forms of the BsubPcs at the F₃BsubPc/Cl-BsubPc heterojunction absorbing at near 615 nm. (bottom) A plot of the log of IPCE versus wavelength for the same device.

band-edge offsets for the frontier orbital energies for $F_5BsubPc$ as an electron donor in heterojunctions with C_{60} as an electron acceptor, and for $F_5BsubPc$ as an electron acceptor in heterojunctions with Cl-BsubPc as the electron donor. This

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was done to evaluate the offsets in transport HOMO and LUMO levels, which are believed to be necessary for efficient exciton dissociation and selective charge extraction.^{4,23,36} Figure 6 shows the offsets in HOMO energies for the $F_5BsubPc/C_{60}$ heterojunction, as estimated from UPS characterization of vacuum deposited thin films of the relevant materials, starting with a clean Au substrate.^{37,38} Equivalent results were obtained with ITO substrates, but electronic equilibrium is easier to maintain during UPS experiments on Au substrates, using film thicknesses of each layer that are the same as those used in the above outlined devices. Above, we used thin layers of vacuum deposited MoO₃ to enhance hole collection in the devices.^{17,18} and this was replicated for the UPS experiments on Au substrates. UPS data is shown in detail in Supporting Information for each heterojunction examined. We typically started with clean Au, added the MoO₃ layer at a device relevant thickness, followed by the donor layer also at its device relevant thickness, and then a few layers of the electron acceptor, up to its final device thickness. After correction for the significant local vacuum level shifts (as determined by shifts in the low kinetic energy region of the UPS spectra), especially for the $F_5BsubPc/C_{60}$ interface, we found an offset in HOMO energies for F₅BsubPc and C₆₀ of 0.6 eV. Assuming a transport bandgap for F₅BsubPc of 2.1 eV and a transport gap of ca. 2.0 eV for C₆₀, we estimate an offset in transport LUMO levels, $E_{\rm F_s}^{\rm LUMO} - E_{\rm C_{60}}^{\rm LUMO}$ = ca. 0.7 eV, which is more than sufficient to provide the driving force for exciton dissociation at the $F_5BsubPc/C_{60}$ interface.²³ There are, in addition, adequate offsets in the HOMO and LUMO levels to provide the rectification and selective charge harvesting associated with planar heterojunction OPVs.³⁹⁻

Figure 7 shows the energetic offsets for the Cl-BsubPc/ F₅BsubPc heterojunctions. The IP of F₅BsubPc is only ca. 0.1 eV larger than for Cl-BsubPc, and from the absorbance spectra of their thin films, the offset in transport LUMOs ($E_{Cl-BsubPc}^{LUMO} - E_{F,BsubPc}^{LUMO}$) is estimated to be ca. 0.1 eV, suggesting a lower driving force for exciton dissociation and offsets in the transport HOMO and LUMO energies, which might make rectification

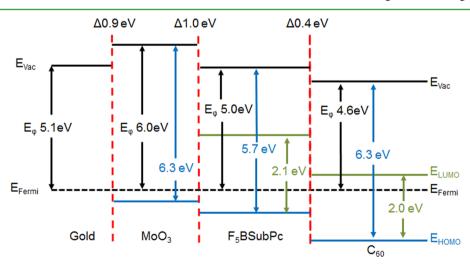


Figure 6. Schematic view of the HOMO levels and local vacuum level shifts for Au/MoO₃/F₅BsubPc/C₆₀ heterojunctions, as determined from layerby-layer vacuum deposition of these materials over a thickness range of ~10 nm (Figure S1, Supporting Information). Electronic equilibrium is maintained with the spectrometer throughout all deposited layers; hence we show the Fermi energy to be constant throughout all layers. The offset in HOMO energies between the F₃BsubPc and C₆₀ layers (0.6 eV) coupled with the assumed bandgap energies for each component (2.1 and 2.0 eV, respectively) leads to an offset in transport LUMO energies ($E_{F_{con}}^{LUMO} - E_{Con}^{LUMO} =$ ca. 0.7 eV), providing for both the rectification seen in these planar heterojunctions in the dark, and the driving force for exciton dissociation under illumination.

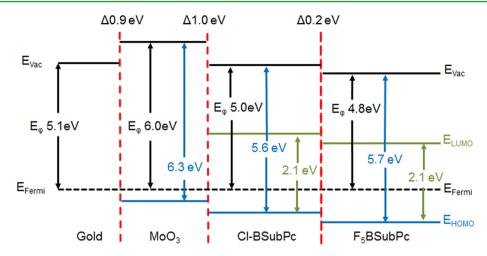


Figure 7. Schematic view of the HOMO levels and local vacuum level shifts for Au/MoO₃/Cl-subPc/F₅BsubPc heterojunctions, as determined from layer-by-layer vacuum deposition of these materials over a thickness range of ~10 nm (Figure S2, Supporting Information). Electronic equilibrium with the spectrometer is maintained throughout all deposited layers, hence the common Fermi energy. The offset in HOMO energies between the Cl-BsubPc and F₅BsubPc (0.1 eV) coupled with the assumed bandgap energies for each component (2.1 eV for each BsubPc) leads to an offset in transport LUMO energies ($E_{Cl}^{LUMO} - E_{F_5}^{LUMO}$ = ca. 0.1 eV).

of the current–voltage response difficult.^{23,24,39} Only small shifts in local vacuum level were seen in the UPS spectra of these heterojunctions (Figure S2, Supporting Information), as revealed by shifts in the low kinetic energy edges of the photoemission spectra. Despite the small offsets and the fact that both Cl-BsubPc and F₅BsubPc show Fermi levels that suggest both materials are p-type (relative to C₆₀), reasonable photovoltaic behavior was observed and detailed above for the Cl-BsubPc/F₅BsubPc heterojunction.

OPVs Based on Multiple Electron Donor or Multiple Electron Acceptor Layers. Cascade OPVs have the potential to make use of three organic materials with monotonically decreasing frontier orbital energy levels.^{43–46} In one example of an energetically monotonic arrangement of materials, the first layer acts an electron donor layer, the second layer is ambipolar (transports both holes and electrons), and the third layer acts as an electron acceptor layer. This configuration may create two heterojunction regions for exciton dissociation (donor/ ambipolar layer and ambipolar layer/acceptor) but the additional heterojunction adds an additional energetic barrier to charge extraction, thus lowering the expected V_{oc} .

As a proof of concept that F5BsubPc could act as the ambipolar component in a cascade arrangement, we explored two device configurations: ITO/MoO₃ (6 nm)/Cl-BsubPc (10 nm)/F₅BsubPc $(x \text{ nm})/C_{60}$ (40 nm)/BCP (10 nm)/Al (100 nm) and ITO/PEN (50 nm)/F₅BsubPc (11 nm)/C₆₀ (40 nm)/BCP (10 nm)/Al (100 nm). For the Cl-BsubPc/ F₅BsubPc heterojunction, the F₅BsubPc layer was varied between 10, 20 and 40 nm and the OPV device performance was evaluated at each thickness. At 20 and 40 nm, an s-kink was observed in the J/V curve, which is likely attributable to enhanced recombination or charge carrier mobility imbalance in the cascade device architecture.²² Higher series resistances were also seen at F5BsubPc thicknesses of 20 and 40 nm. However, at a F₅BsubPc layer thickness of 10 nm, power conversion efficiency of 0.96% and a standard shaped J/V curve was observed (Figure 8). By comparison to the planar bilayer devices described above, this cascade device produces ~50% more photocurrent but has a V_{oc} smaller by ca. 0.3 V (Table 3).

PEN/F₅BsubPc/C₆₀ cascade devices also showed the same general performance; higher photocurrents with lower V_{oc} 's.

However, in this case less significant gains are observed in overall cell performance perhaps since the PEN/F₅BsubPc/C₆₀ cascade device combines two sub cells with strikingly mismatched device performance (V_{oc} of 0.71 V and a J_{sc} of 1.43 mA/cm², Table 3). Nonetheless, given these two examples, we can conclude that F₅BsubPc can also function as an ambipolar layer in a cascade arrangement, transporting both electrons and holes while also absorbing light.

CONCLUSIONS

We have demonstrated that F₅BsubPc can function as either an electron-donating/hole-transporting material or an electronaccepting/electron-transporting material in a planar heterojunction organic photovoltaic (PHJ OPV) configuration. Using C₆₀ as the electron acceptor and F₅BsubPc as the electron donor, good frontier orbital energy offsets were observed, which are consistent with other high IP donor/C₆₀ heterojunctions. OPVs with good rectification, high V_{oc} 's, and reasonable efficiencies were achieved for unoptimized devices. For Cl-BsubPc/F5BsubPc devices, the offsets in both the HOMO and LUMO energies were measured to be small (ca. 0.1 eV), but good rectification and photovoltaic activity was still achieved. In each case, by deconstructing the devices, we could verify that the rectifying and photocurrent producing interfaces were the F₅BsubPc/C₆₀ and Cl-BsubPc/F₅BsubPc interfaces. In each of the device configurations, a strong contribution to the photocurrent was observed at energies below the lowest energy excitonic state. Charge transfer complexes at the donor/ acceptor interface are observed by IPCE and are ~ 1.7 eV for both F₅BsubPc/C₆₀ and Cl-BsubPc/F₅BsubPc.

To probe the ambipolarity of $F_{s}BsubPc$, we also constructed cascade OPV devices. We demonstrate that $F_{s}BsubPc$ can function in this role by effectively transporting both holes and electrons while absorbing light albeit only for a specific (and largely unoptimized) $F_{s}BsubPc$ layer thickness. Our cascade devices produced, on average, higher photocurrent (~50%) presumably due to the two exciton dissociation interfaces but a reduced open circuit voltage (by comparison to a similar donor/acceptor OPV).^{47,48}

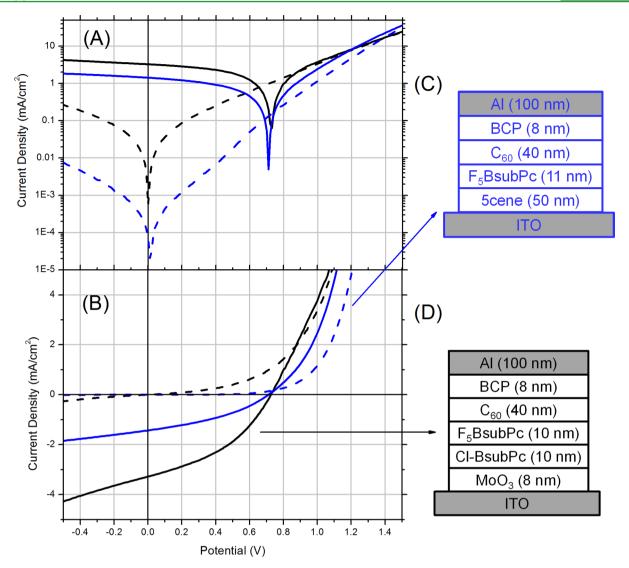


Figure 8. Log (A) and linear (B) dark and illuminated J/V plots for OPV cascade devices in configurations that require $F_5BsubPc$ to form a heterojunction with PEN backed with C_{60} as the electron transport layer (blue line) or a heterojunction with Cl-BsubPc (black line) backed with C_{60} as the electron transport layer. (C) Device configuration of PEN/ $F_5BsubPc/C_{60}$ OPV. (D) Device configuration of Cl-BsubPc/ $F_5BsubPc/C_{60}$ OPV.

Table 3. A Comparison between the Cascade OPV Devices ITO/MoO₃/Cl-BsubPc/F₅BsubPc/C₆₀/BCP/Al and ITO/PEN/F₅BsubPc/C₆₀/BCP/Al and Their Comparative Bilayer Devices

donor (nm)	F ₅ BsubPc (nm)	C ₆₀ (nm)	$V_{ m oc} \ ({ m V})$	$J_{\rm sc}$ (mA/cm ²)	FF (%)	η (%)
Cl-BsubPc (10)	10	40	0.73	3.28	40.3	0.97
	17	40	1.02	2.03	46.2	0.96
Cl-BsubPc (10)	20		0.98	1.72	44.9	0.76
PEN (50)	29		0.82	0.46	40.8	0.15
PEN (50)	11	40	0.71	1.43	37.2	0.38

 $F_5BsubPc$ is a representative of the more general class of phenoxy-BsubPcs. We feel these results further the idea that the field should look beyond the prototypical BsubPc, Cl-BsubPc, and towards the large synthetic variation that is available to phenoxy-BsubPcs. Based on the work outlined herein and our previous work showing the dual functionality of $F_5BsubPc$ in OLEDs,^{6,7} we feel phenoxy-BsubPcs may have significant

promise across multiple organic electronic device configurations.

EXPERIMENTAL SECTION

ITO coated glass was obtained from Colorado Concept Coating LLC., with a thickness of 150 nm and a sheet resistance of 15 Ω per square (four-point probe). Substrates (1 in. \times 1 in.) were cleaned by (1) washing with nanopure water and scrubbing with 10% Triton X-100, (2) sonication in 10% Triton X-100 for 15 min, (3) rinsing and scrubbing with nanopure water, (4) sonication in nanopure water for 15 min, (5) rinsing and scrubbing with 100% ethanol, and (6) sonication in 200 proof ethanol for 15 min. Large area (0.125 cm²) substrates of 1 in. \times 1 in. size were rinsed with 100% ethanol and then patterned with S1813 positive photoresist (Rohm and Haas), followed by exposure, development, etching, and removal of the photoresist to create the necessary patterns for 0.125 cm² substrates. Each patterned ITO substrate was etched using aqua regia at a 3:1 ratio at 120 °C for ca. 30 s. The substrates were then cleaned using the six step procedure above and stored in 100% ethanol.

Prior to introduction to the vacuum systems, cleaned ITO substrates were removed from the ethanol solution and dried under a stream of nitrogen. For the plasma treated experiments, the etched substrates were put into a Harrick PDC-32G plasma generator, which

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was at ca. 0.4 Torr with dry oxygen etching occurring at 10 min. Detergent solvent cleaned (DSC) ITO was used directly from the stored substrates in ethanol. C₆₀, copper phthalocyanine (CuPc), and pentacene (Mer Corp, Aldrich) were triply sublimed using an in house sublimation system, and bathocuproine (Aldrich) was sublimed twice. Molybdenum oxide (99.99%) was purchased from Aldrich as used as is. The synthesis and purification of pentafluorophenoxy boron subphthalocyanine (F₅BsubPc) has been previously described.⁵ In order to ensure high purity, F5BsubPc was doubly sublimed in a thermal gradient sublimation apparatus. A custom high vacuum deposition system was used to fabricate the various organic photovoltaic test cells used in this study, as described previously.⁴⁹⁻⁵¹ All molecules were thoroughly degassed in the vacuum system prior to use. Aluminum top contacts (Alfa Aesar) were deposited with 99.999% rated material. All molecules used in vacuum deposition were loaded into a Knudsen cell type system and deposited at pressures equal to or lower than 9×10^{-7} Torr. Thin films were deposited at a rate of approximately 0.5 Å/s for organics and 1 Å/s for aluminum, measured with a 10 MHz quartz crystal microbalance (QCM-Newark) and a frequency monitor (Agilent, model 53131A) or a 6 MHz QCM (Lesker) monitored via an Inficon Q-pod QCM monitor. Top contacts were deposited via thermal evaporation and measured with a 6 MHz QCM (Tangydine) and a frequency monitor (Inficon, model 758-500-G1). The OPV devices fabricated in this study were each constructed with pentafluorophenoxy boron subphthalocyanine (F₅BsubPc) as an active material in a variety of planar heterojunction OPVs. F₅BsubPc was paired with other common organic semiconductors so as to limit the number of unknown factors in this investigation. Thicknesses of the active layers were optimized by finding those thicknesses that produced maximum photocurrent and photopotentials and minimized series resistance, maximizing fillfactors. In each case, device performance was reported as an average of more than eight circular 0.019 cm² cells. The devices were constructed on a transparent ITO glass electrode with a vacuum deposited 100 nm layer of Al as a top electrode. In each cell, a layer of BCP was used as a barrier to Al migration, which also prevents or reduces the amount of pin-holes and acts as an electron selective interface layer.⁵² Cell configurations were either $ITO/F_5BsubPc/C_{60}/BCP/Al$, ITO/Cl-BsubPc/F5BsubPc/BCP/Al, or ITO/PEN/F5BsubPc/BCP/Al.

X-ray photoelectron spectroscopy (XPS) and UV-photoelectron spectroscopy (UPS) were used to probe the interface between each material. Aluminum K α X-rays at 1486.6 eV were used for XPS measurements, and He(I) was used as the source (21.2 eV) for UPS measurements (Kratos Axis-Ultra spectrophotometer). For UPS measurements, the sample was biased at 10 V to enhance the yield of low kinetic energy electrons.^{36,38,53,54} The spot size was 300 μ m by 700 μ m for XPS and 1 mm to 3 mm for UPS measurements. The work function of clean Au (Alfa Aesar) was verified to be ca. 5.1 ± 0.1 eV and to set the kinetic energy of photoelectrons emitted at the Fermi energy for Au, which was assumed constant for all subsequent ITO/ organic heterojunctions; that is, we assume electronic equilibrium between the sample and the spectrometer.³⁶ Throughout the data collection process, photoemission from the gold Fermi level varied from 30.7 to 34 eV (absolute kinetic energy) for sputter cleaned gold substrates. Gold was polished to a mirror finish and soaked in piranha for 24 h and subsequently argon sputtered prior to the experiment. Molecules were deposited on ITO (Colorado Concept Coating LLC) in an attached high vacuum deposition system where the rate was monitored using a quartz crystal monitor (Agilent 53131A) and custom designed quartz oscillator (QCM-Newark) where depositions of sub-angstrom per second are possible on either gold or ITO substrates. The temperature was controlled by an Omega CN76000 temperature driver/controller.

For testing devices, a custom made current–voltage (J/V) system was used to test six devices per substrate for large area (0.125 cm^2) devices and 25 per substrate for small area (0.019 cm^2) . J/Vmeasurements were performed in a nitrogen filled glovebox (Mbraun Labmaster) where water and oxygen levels were below 0.1 ppm. A 300 W xenon arc lamp (Newport) that is current controlled was used as the light source. The beam path from the source to the sample was filtered using an AM 1.5 filter (MellesGriot) to simulate the solar spectrum. The filtered light was then optically diffused using an engineered diffuser with an output of 40° (Newport). The power density at the surface of the devices was tuned with a flat response thermopile (Newport) and cross checked with a calibrated silicon photodiode (Newport, model 818-SL with OD3 Attenuator) to get 100 mW/cm^2 . The potential applied to the device was swept using a source meter (Keithley 2400) and in-house software (National Instruments Labview 8.2). The potential was swept from -1.00 to 1.50 V using a 5.00 mV step starting from the negative potential. Absorbance measurements were carried out with a UV/visible spectrophotometer (Agilent Technologies, model 8453). Incident and absorbed power conversion efficiency (IPCE and APCE) measurements were acquired in an in-house fabricated hermetically sealed unit. A 300 W xenon arc lamp (Newport) was used as the light source. The 256 Hz modulated light (Stanford Research) was passed through an automated monochromator (Newport, Cornerstone 130), where the data was acquired at 4 nm intervals. Incident power though the monochromator and custom optics setup was measured with a silicon photodiode (Newport model 818-SL), while device measurements were acquired at short-circuit. The resulting signal was fed to a lock-in amplifier (EG&G model 5209), while the output was fed into an in-house data acquisition system (National Instruments Labview 8.2). The APCE values have been normalized because refractive indices and reflection off the back electrode were not taken into account.

ASSOCIATED CONTENT

S Supporting Information

Full ultraviolet photoelectron spectroscopy data of the $F_5BsubPc/C_{60}$ and Cl-BsubPc/ $F_5BsubPc$ interface, large area device performance of the Cl-BsubPc/ $F_5BsubPc$ device, and Cl-BsubPc Schottky device performance. This material is available free of charge via the Internet at http://pubs.acs.org.

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Notes

The authors declare no competing financial interest.

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REFERENCES

(1) Shun-Wei, L.; Chin-Ti, C.; Jiun-Haw, L. Presented at the Nanoelectronics Conference (INEC), 2011 IEEE 4th International, 21-24 June 2011; 2011, 1-2.

(2) Liu, S.-W.; Lin, C.-F.; Lee, C.-C.; Su, W.-C.; Chen, C.-T.; Lee, J.-H. J. Electrochem. Soc. **2012**, 159 (2), H191–H194.

- (3) Jiang, H. Macromol. Rapid Commun. 2010, 31 (23), 2007–2034.
 (4) Mutolo, K. L.; Mayo, E. I.; Rand, B. P.; Forrest, S. R.; Thompson,
- M. E. J. Am. Chem. Soc. 2006, 128 (25), 8108–8109.

(5) Gommans, H.; Aernouts, T.; Verreet, B.; Heremans, P.; Medina, A.; Claessens, C. G.; Torres, T. *Adv. Funct. Mater.* **2009**, *19* (21), 3435–3439.

ACS Applied Materials & Interfaces

(6) Sullivan, P.; Duraud, A.; Hancox, I.; Beaumont, N.; Mirri, G.; Tucker, J. H. R.; Hatton, R. A.; Shipman, M.; Jones, T. S. *Adv. Energy Mater.* **2011**, *1* (3), 352–355.

- (7) Morse, G. E.; Helander, M. G.; Maka, J. F.; Lu, Z. H.; Bender, T. P. ACS Appl. Mater. Interfaces **2010**, 2 (7), 1934–1944.
- (8) Helander, M. G.; Morse, G. E.; Qiu, J.; Castrucci, J. S.; Bender, T. P.; Lu, Z. H. ACS Appl. Mater. Interfaces **2010**, 2 (11), 3147–3152.
- (9) Díaz, D. D.; Bolink, H. J.; Cappelli, L.; Claessens, C. G.; Coronado, E.; Torres, T. *Tetrahedron Lett.* **2007**, *48* (27), 4657–4660.
- (10) Yasuda, T.; Tsutsui, T. Mol. Cryst. Liq. Cryst. **2006**, 462 (1), 3–9.
- (11) Castrucci, J. S.; Helander, M. G.; Morse, G. E.; Lu, Z. H.; Yip, C. M.; Bender, T. P. *Cryst. Growth Des.* **2012**, *12* (3), 1095–1100.
- (12) Morse, G. E.; Castrucci, J. S.; Helander, M. G.; Lu, Z. H.; Bender, T. P. ACS Appl. Mater. Interfaces **2011**, 3 (9), 3538–3544.
- (13) Beaumont, N.; Cho, S. W.; Sullivan, P.; Newby, D.; Smith, K. E.; Jones, T. S. Adv. Funct. Mater. **2012**, 22, 561–566.
- (14) Morse, G. E.; Bender, T. P. ACS Appl. Mater. Interfaces 2012, 4 (10), 5055-5068.
- (15) Gommans, H.; Cheyns, D.; Aernouts, T.; Girotto, C.; Poortmans, J.; Heremans, P. *Adv. Funct. Mater.* **2007**, *17* (15), 2653–2658.
- (16) Hancox, I.; Sullivan, P.; Chauhan, K.; Beaumont, N.; Rochford, L.; Hatton, R.; Jones, T. Org. Electron. **2010**, *11* (12), 2019–2025.
- (17) Kroger, M.; Hamwi, S.; Meyer, J.; Riedl, T.; Kowalsky, W.; Kahn, A. Appl. Phys. Lett. 2009, 95 (12), 3.
- (18) Kroger, M.; Hamwi, S.; Meyer, J.; Riedl, T.; Kowalsky, W.; Kahn, A. Org. Electron. 2009, 10 (5), 932–938.
- (19) Hamwi, S.; Meyer, J.; Kröger, M.; Winkler, T.; Witte, M.; Riedl, T.; Kahn, A.; Kowalsky, W. *Adv. Funct. Mater.* **2010**, 20 (11), 1762–1766.
- (20) Tress, W.; Petrich, A.; Hummert, M.; Hein, M.; Leo, K.; Riede, M. *Appl. Phys. Lett.* **2011**, *98*, No. 063301.
- (21) Tress, W.; Leo, K.; Riede, M. Phys. Rev. B 2012, 85, No. 155201.
 (22) Mandoc, M. M.; Koster, L. J. A.; Blom, P. W. M. Appl. Phys. Lett.
- **2007**, *90*, No. 133504.
- (23) Rand, B. P.; Burk, D. P.; Forrest, S. R. Phys. Rev. B 2007, 75 (11), No. 115327.
- (24) Forrest, S. R. MRS Bull. 2005, 30 (1), 28-32.
- (25) Potscavage, W. J.; Yoo, S.; Kippelen, B. *Appl. Phys. Lett.* 2008, 93, No. 193308.
- (26) Zhang, M.; Ding, H.; Gao, Y.; Tang, C. Appl. Phys. Lett. 2010, 96 (18), No. 183301.
- (27) Shrotriya, V.; Li, G.; Yao, Y.; Moriarty, T.; Emery, K.; Yang, Y. *Adv. Funct. Mater.* **2006**, *16* (15), 2016–2023.
- (28) Pandey, R.; Holmes, R. J. Adv. Mater. 2010, 22 (46), 5301-5305.
- (29) Gommans, H.; Verreet, B.; Rand, B. P.; Muller, R.; Poortmans, J.; Heremans, P.; Genoe, J. Adv. Funct. Mater. 2008, 18 (22), 3686–3691.
- (30) Graham, K. R.; Erwin, P.; Nordlund, D.; Vandewal, K.; Li, R.; Ngongang Ndjawa, G. O.; Hoke, E. T.; Salleo, A.; Thompson, M. E.; McGehee, M. D.; Amassian, A. *Adv. Mater.* **2013**, *25*, 6076–6082.
- (31) Verreet, B.; Rand, B. P.; Cheyns, D.; Hadipour, A.; Aernouts, T.; Heremans, P.; Medina, A.; Claessens, C. G.; Torres, T. *Adv. Energy Mater.* **2011**, *1* (4), 565–568.
- (32) Griffith, O. L.; Anthony, J. E.; Jones, A. G.; Lichtenberger, D. L. J. Am. Chem. Soc. **2009**, 132 (2), 580–586.
- (33) Jurchescu, O. D.; Baas, J.; Palstra, T. T. M. Appl. Phys. Lett. **2004**, 84 (16), 3061–3063.
- (34) Yoo, S.; Potscavage, W. J., Jr.; Domercq, B.; Han, S.-H.; Li, T.-D.; Jones, S. C.; Szoszkiewicz, R.; Levi, D.; Riedo, E.; Marder, S. R.; Kippelen, B. *Solid State Electron.* **2007**, *51* (10), 1367–1375.
- (35) Pakhomov, G.; Travkin, V. V.; Bogdanova, A. Y.; Guo, T. F. J. Porphyrins Phthalocyanines **2008**, *12* (11), 1182–1186.
- (36) Armstrong, N. R.; Wang, W. N.; Alloway, D. M.; Placencia, D.; Ratcliff, E.; Brumbach, M. *Macromol. Rapid Commun.* **2009**, 30 (9-10), 717–731.

- (37) Ishii, H.; Sugiyama, K.; Ito, E.; Seki, K. *Adv. Mater.* **1999**, *11* (8), 605–625.
- (38) Hwang, J.; Wan, A.; Kahn, A. Mater. Sci. Eng., R: Reports 2009, 64 (1-2), 1-31. (c) Cahen, D.; Kahn, A. Adv. Mater. 2003, 15 (4), 271-277
- (39) Tang, C. W. Appl. Phys. Lett. 1986, 48 (2), 183-185.
- (40) Zhang, M. L.; Wang, H.; Tang, C. W. Appl. Phys. Lett. 2010, 97 (14), No. 143503.
- (41) Alloway, D. M.; Armstrong, N. R. Appl. Phys. A: Mater. Sci. Process. 2009, 95 (1), 209–218.
- (42) Ratcliff, E. L.; Zacher, B.; Armstrong, N. R. J. Phys. Chem. Lett. 2011, 1337-1350.
- (43) Zhang, G.; Li, W.; Chu, B.; Chen, L.; Yan, F.; Zhu, J.; Chen, Y.; Lee, C. S. Appl. Phys. Lett. **2009**, 94 (14), No. 143302.
- (44) Ichikawa, M.; Takekawa, D.; Jeon, H. G.; Banoukepa, G. D. R. Org. Electron. 2013, 14 (3), 814-820.
- (45) Schlenker, C. W.; Barlier, V. S.; Chin, S. W.; Whited, M. T.; McAnally, R. E.; Forrest, S. R.; Thompson, M. E. *Chem. Mater.* 2011, 23 (18), 4132–4140.
- (46) Chen, M. C.; Liaw, D. J.; Huang, Y. C.; Wu, H. Y.; Tai, Y. Sol. Energy Mater. Sol. Cells 2011, 95 (9), 2621–2627.
- (47) Yuen, A. P.; Hor, A.-M.; Preston, J. S.; Klenkler, R.; Bamsey, N. M.; Loutfy, R. O. Appl. Phys. Lett. **2011**, 98 (17), No. 173301.
- (48) Huang, J.-H.; Velusamy, M.; Ho, K.-C.; Lin, J.-T.; Chu, C.-W. J. Mater. Chem. 2010, 20 (14), 2820–2825.
- (49) Wang, W.; Placencia, D.; Armstrong, N. R. Org. Electron. 2011, 12, 383-393.
- (50) Placencia, D.; Wang, W. N.; Gantz, J.; Jenkins, J. L.; Armstrong, N. R. J. Phys. Chem. C 2011, 115 (38), 18873–18884.
- (51) Placencia, D.; Wang, W. N.; Shallcross, R. C.; Nebesny, K. W.; Brumbach, M.; Armstrong, N. R. *Adv. Funct. Mater.* **2009**, *19* (12), 1913–1921.
- (52) Wang, Y.-M.; Teng, F.; Zhou, Q.-C.; Wang, Y.-S. Appl. Surf. Sci. 2006, 252 (6), 2355–2359.
- (53) Schlaf, R.; Parkinson, B.; Lee, P.; Nebesny, K.; Armstrong, N. J. Phys. Chem. B **1999**, 103 (15), 2984–2992.
- (54) Alloway, D. M.; Graham, A. L.; Yang, X.; Mudalige, A.; Colorado, R.; Wysocki, V. H.; Pemberton, J. E.; Randall Lee, T.; Wysocki, R. J.; Armstrong, N. R. J. Phys. Chem. C **2009**, 113 (47), 20328–20334.